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# Symplectic volumes of double weight varieties associated with $SU(3)$ , I

By

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## Abstract

We consider double weight varieties, that is, symplectic torus quotients for a direct product of two integral coadjoint orbits of  $SU(3)$ , and investigate their symplectic volumes. According to a fundamental theorem for symplectic quotients, it is equivalent to studying weight multiplicities in a tensor product of two irreducible representations of  $SU(3)$ , and their asymptotic behavior. We assume that both of two coadjoint orbits used to define the double weight variety are flag manifolds of  $SU(3)$ . As a main result, we obtain an explicit formula for the symplectic volumes of double weight varieties.

## § 1. Introduction

Let  $M$  be a symplectic manifold with a Hamiltonian action of a torus  $T$  and  $\Phi : M \rightarrow \mathfrak{t}^*$  its moment map. For any regular value  $\mu \in \mathfrak{t}^*$  of  $\Phi$ , the symplectic quotient at  $\mu$  is defined by

$$M//_{\mu}T := \Phi^{-1}(\mu)/T.$$

For example, let  $M$  be a coadjoint orbit  $\mathcal{O}_{\lambda}$  of a compact semi-simple Lie group  $G$  through the point  $\lambda \in \mathfrak{t}^*$ , and consider the action of the maximal torus  $T \subset G$  on  $\mathcal{O}_{\lambda}$ . Then  $M//_{\mu}T = \mathcal{O}_{\lambda}//_{\mu}T$  is called a *weight variety*. Many results have been obtained about weight varieties. Knutson [15] studied the relation with weight spaces of an irreducible representation (thereby  $\mathcal{O}_{\lambda}//_{\mu}T$  were named weight varieties), Guillemin-Lerman-Sternberg [6] gave some formulas for the volumes of weight varieties, and Goldin [4] gave an explicit expression for the cohomology rings of weight varieties for

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$G = SU(n)$ . For compact semi-simple Lie groups except of type  $A$ , it is known that weight varieties are orbifolds in general (see [5]).

For  $\lambda_1, \lambda_2 \in \mathfrak{t}^*$ , let  $M$  be a direct product  $\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}$  of two coadjoint orbits of  $G$ , and consider the diagonal action of  $T \subset G$  on  $\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}$ . Then  $M//_{\mu}T = (\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T$  is called a *double weight variety*.

In this paper, we consider the case where  $G = SU(3)$ . Except for the orbit consisting only of the origin, each coadjoint orbit of  $SU(3)$  is diffeomorphic to either the flag manifold  $SU(3)/T$  or the complex projective space  $\mathbb{P}^2$ . Our aim is to express the symplectic volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T)$  of the double weight variety in an explicit form.

First, we assume that  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$ , where  $P$  denotes the weight lattice of  $G$  and  $P_+$  denotes the set of dominant integral weights of  $G$ . As we will discuss in Section 2, under certain conditions on  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$ , we can express the volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T)$  in terms of representations of  $T$ . Namely,  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T)$  is equal to

$$\lim_{k \rightarrow \infty} \frac{1}{k^d} \cdot [V_{k\lambda_1} \otimes V_{k\lambda_2} : W_{k\mu}],$$

where  $V_{\lambda}$  denotes the irreducible representation of  $G$  with highest weight  $\lambda \in P_+$ , and for a representation  $V$  of  $T$ ,  $[V : W_{\mu}]$  denotes the multiplicity of weight space  $W_{\mu}$  ( $\mu \in P$ ) in the weight decomposition of  $V$ . Moreover,  $k$  runs over positive integers and  $d$  is the complex dimension of  $(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T$ . Here, the theorem of Guillemin-Sternberg (and its generalization) on the characteristic numbers of symplectic quotients (see, e.g., [8] and [17]) plays the key role, as well as the Borel-Weil theorem and the Hirzebruch-Riemann-Roch theorem. The argument above are essentially the same with those in [20] and [21].

Next, we assume that both of two coadjoint orbits  $\mathcal{O}_{\lambda_1}$  and  $\mathcal{O}_{\lambda_2}$  are flag manifolds of  $SU(3)$ . Namely, we assume  $\lambda_1, \lambda_2 \in P_{++}$ , where  $P_{++}$  denotes the set of dominant integral weights which belong to the interior of the Weyl chamber. The main result in this paper is an explicit formula for the volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T)$ . The details of the notation will be given in Sections 2 and 3.

**Theorem.** (See Theorem 3.4 below.) *Let  $\lambda_1, \lambda_2 \in P_{++}$  and  $\mu \in P$  satisfy the assumptions (A0), (A1) and (A2) in Section 3.2. Then we have*

$$\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T) = \sum_{w_1, w_2 \in W} \varepsilon(w_1)\varepsilon(w_2)F(\lambda_1, \lambda_2, \mu; w_1, w_2),$$

where  $W$  denotes the Weyl group of  $SU(3)$  and  $F(\lambda_1, \lambda_2, \mu; w_1, w_2)$  is given as follows.

(1) If  $w_1\lambda_1 + w_2\lambda_2 - \mu \in \gamma_1$ ,

$$F(\lambda_1, \lambda_2, \mu; w_1, w_2) = \frac{1}{12} \langle w_1\lambda_1 + w_2\lambda_2 - \mu, \Lambda_2 \rangle^3 \langle w_1\lambda_1 + w_2\lambda_2 - \mu, 2\Lambda_1 - \Lambda_2 \rangle.$$

(2) If  $w_1\lambda_1 + w_2\lambda_2 - \mu \in \gamma_2$ ,

$$F(\lambda_1, \lambda_2, \mu; w_1, w_2) = \frac{1}{12} \langle w_1\lambda_1 + w_2\lambda_2 - \mu, \Lambda_1 \rangle^3 \langle w_1\lambda_1 + w_2\lambda_2 - \mu, -\Lambda_1 + 2\Lambda_2 \rangle.$$

(3) Otherwise,

$$F(\lambda_1, \lambda_2, \mu; w_1, w_2) = 0.$$

## § 2. Preliminaries

### § 2.1. The representation theory of $SU(3)$ and notation

We review some standard facts about the representation theory of  $SU(3)$  in order to fix our notation. We refer to [2] for the generalities on compact Lie groups and their representations.

Let  $G = SU(3)$ ,  $\mathfrak{g} = \mathfrak{su}(3)$ ,  $T$  the standard maximal torus of  $G$  consisting of diagonal matrices in  $G$ , and  $\mathfrak{t}$  its Lie algebra. Let  $\mathfrak{g}^*$  and  $\mathfrak{t}^*$  be the duals of  $\mathfrak{g}$  and  $\mathfrak{t}$ , respectively. We denote by  $\langle \cdot, \cdot \rangle$  the pairing between  $\mathfrak{g}^*$  and  $\mathfrak{g}$ , or  $\mathfrak{t}^*$  and  $\mathfrak{t}$ . Let  $W \cong \mathfrak{S}_3$  be the Weyl group of  $G = SU(3)$  with respect to  $T$ . We define an  $\text{Ad}G$ -invariant positive definite inner product  $(\cdot, \cdot)$  on  $\mathfrak{g}$  by

$$(X, Y) := -\frac{1}{4\pi^2} \text{Tr}(XY) \quad (X, Y \in \mathfrak{g}).$$

We identify  $\mathfrak{g}^*$  with  $\mathfrak{g}$  by the inner product  $(\cdot, \cdot)$ . We regard  $\mathfrak{t}^*$  as a subspace of  $\mathfrak{g}^*$  by the identification

$$\mathfrak{t}^* = \{f \in \mathfrak{g}^* | t \cdot f = f \ (\forall t \in T)\}.$$

The elements

$$(2.1) \quad H_1 = 2\pi\sqrt{-1} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad H_2 = 2\pi\sqrt{-1} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

in  $\mathfrak{t}$  are generators of the integral lattice  $\text{Ker}(\exp : \mathfrak{t} \rightarrow T)$  and form a basis of  $\mathfrak{t}$ . Under the identification by the inner product, we define the simple roots  $\alpha_1, \alpha_2$  in  $\mathfrak{t}^*$  by the elements which correspond to  $H_1, H_2 \in \mathfrak{t}$ , respectively. Let  $\Delta_+ := \{\alpha_1, \alpha_2, \alpha_1 + \alpha_2\}$  be the positive root system. Let us set

$$Q := \mathbb{Z}\alpha_1 + \mathbb{Z}\alpha_2,$$

$$\gamma_1 := \mathbb{R}_{>0}\alpha_1 + \mathbb{R}_{>0}(\alpha_1 + \alpha_2), \quad \gamma_2 := \mathbb{R}_{>0}(\alpha_1 + \alpha_2) + \mathbb{R}_{>0}\alpha_2,$$

$$\bar{\gamma}_1 := \mathbb{R}_{\geq 0}\alpha_1 + \mathbb{R}_{\geq 0}(\alpha_1 + \alpha_2), \quad \bar{\gamma}_2 := \mathbb{R}_{\geq 0}(\alpha_1 + \alpha_2) + \mathbb{R}_{\geq 0}\alpha_2.$$

We define the fundamental weights  $\Lambda_1, \Lambda_2 \in \mathfrak{t}^*$  by

$$\Lambda_1 := \frac{2\alpha_1 + \alpha_2}{3}, \Lambda_2 := \frac{\alpha_1 + 2\alpha_2}{3}.$$

Under the identification by the inner product,  $\Lambda_1, \Lambda_2 \in \mathfrak{t}^*$  correspond to the elements

$$I_1 = \frac{2\pi\sqrt{-1}}{3} \begin{pmatrix} 2 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \quad I_2 = \frac{2\pi\sqrt{-1}}{3} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

in  $\mathfrak{t}$ , respectively. Define

$$\begin{aligned} \mathfrak{t}_+^* &:= \mathbb{R}_{\geq 0}\Lambda_1 + \mathbb{R}_{\geq 0}\Lambda_2, \quad \mathfrak{t}_{++}^* := \mathbb{R}_{> 0}\Lambda_1 + \mathbb{R}_{> 0}\Lambda_2, \\ P &:= \mathbb{Z}\Lambda_1 + \mathbb{Z}\Lambda_2, \quad P_+ := \mathbb{Z}_{\geq 0}\Lambda_1 + \mathbb{Z}_{\geq 0}\Lambda_2, \quad P_{++} := \mathbb{Z}_{> 0}\Lambda_1 + \mathbb{Z}_{> 0}\Lambda_2. \end{aligned}$$

The set  $\mathfrak{t}_+^*$  is a fundamental domain of the action of the Weyl group  $W$  on  $\mathfrak{t}^*$ . Let us set

$$\rho := \frac{1}{2} \sum_{\alpha \in \Delta_+} \alpha = \Lambda_1 + \Lambda_2.$$

According to the representation theory of compact Lie groups, irreducible representations of  $G$  are, by assigning their highest weights, in one-to-one correspondence with elements in  $P_+$ . For  $\lambda \in P_+$ , we denote by  $V_\lambda$  the irreducible representation of  $G$  with the highest weight  $\lambda \in P_+$ , and by  $\chi_\lambda : G \rightarrow \mathbb{C}$  the character of  $V_\lambda$ . For  $\mu \in P$ , we define  $e^\mu : T \rightarrow \mathbb{C}$  by  $e^\mu(t) := e^{2\pi\sqrt{-1}\langle \mu, X \rangle}$  for  $t = \exp X \in T$  ( $X \in \mathfrak{t}$ ). By the Weyl character formula,  $\chi_\lambda$  is given as a function on  $T$  by

$$\chi_\lambda(t) = \frac{\sum_{w \in W} \varepsilon(w) e^{w(\lambda + \rho)}(t)}{e^\rho(t) \prod_{\alpha \in \Delta_+} (1 - e^{-\alpha}(t))},$$

where  $\varepsilon(w) = \pm 1$  is the signature of  $w \in W$ .

## § 2.2. Coadjoint orbits

Although we mainly consider the case  $G = SU(3)$ , most of the following still holds when  $G$  is a general compact Lie group. For further details on coadjoint orbits, see, e.g., [13] and [16]. We also refer to [12] for the Borel-Weil theorem.

The left coadjoint action of  $G$  on  $\mathfrak{g}^*$  is defined by  $g \cdot f := \text{Ad}^*(g)f$  for  $g \in G$  and  $f \in \mathfrak{g}^*$ , where

$$\langle \text{Ad}^*(g)f, X \rangle = \langle f, \text{Ad}(g^{-1})X \rangle \quad (X \in \mathfrak{g}).$$

We denote by  $\mathcal{O}_\lambda = G \cdot \lambda$  the coadjoint orbit through  $\lambda \in \mathfrak{t}^* \subset \mathfrak{g}^*$ . Then the intersection  $\mathcal{O}_\lambda \cap \mathfrak{t}^*$  is the  $W$ -orbit through  $\lambda$ , and  $\mathcal{O}_\lambda \cap \mathfrak{t}_+^*$  consists of a single point. In other words, coadjoint orbits are parametrized by elements in  $\mathfrak{t}_+^*$ .

In particular, for  $G = SU(3)$ , coadjoint orbits  $\mathcal{O}_\lambda$  are classified as follows, where  $G_\lambda$  denotes the isotropy subgroup at  $\lambda \in \mathfrak{t}_+^*$  for the coadjoint action of  $G$  on  $\mathfrak{g}^*$ .

- (1) If  $\lambda \in \mathfrak{t}_{++}^*$ , then  $G_\lambda \cong T$  and  $\mathcal{O}_\lambda \cong SU(3)/T$ .
- (2) If  $\lambda \in \mathfrak{t}_+^* - \mathfrak{t}_{++}^* - \{0\}$ , then  $G_\lambda \cong S(U(1) \times U(2))$  and  $\mathcal{O}_\lambda \cong \mathbb{P}^2$ .
- (3) If  $\lambda = 0$ , then  $G_\lambda = G$  and  $\mathcal{O}_\lambda = \{0\}$ .

On each coadjoint orbit  $\mathcal{O}_\lambda$ , there exists a natural  $G$ -invariant symplectic structure  $\omega_\lambda$ , called the Kirillov-Kostant-Souriau symplectic form, defined by

$$(\omega_\lambda)_x(\tilde{X}_x, \tilde{Y}_x) = \langle x, [X, Y] \rangle \quad (x \in \mathcal{O}_\lambda, X, Y \in \mathfrak{g}),$$

where  $\tilde{X}$  is the vector field on  $\mathcal{O}_\lambda$  given by

$$\tilde{X}_x := \left. \frac{d}{dt} \right|_{t=0} (\exp tX) \cdot x.$$

The action of  $G$  on  $\mathcal{O}_\lambda$  is Hamiltonian and the associated moment map is given by the inclusion  $\iota_\lambda : \mathcal{O}_\lambda \hookrightarrow \mathfrak{g}^*$ , that is, we have  $d\langle \iota_\lambda, X \rangle(\cdot) = \omega_\lambda(\tilde{X}, \cdot)$ .

In addition, there exists a  $G$ -invariant complex structure  $J_\lambda$  on  $\mathcal{O}_\lambda$ , which is compatible with the symplectic structure  $\omega_\lambda$ , that is,  $\omega_\lambda(\cdot, J_\lambda \cdot)$  becomes a Riemannian metric on  $\mathcal{O}_\lambda$ , and makes  $\mathcal{O}_\lambda$  into a Kähler manifold.

Moreover, in the case  $\lambda \in P_+$ , there exists a  $G$ -equivariant holomorphic line bundle  $L_\lambda$  over  $\mathcal{O}_\lambda$  such that  $c_1(L_\lambda) = [\omega_\lambda]$ . The Borel-Weil theorem shows that

$$H^0(\mathcal{O}_\lambda, \mathcal{O}(L_\lambda)) \cong V_\lambda, \quad H^k(\mathcal{O}_\lambda, \mathcal{O}(L_\lambda)) = 0 \quad (k > 0)$$

as representations of  $G$ , where  $\mathcal{O}(L_\lambda)$  denotes the sheaf of germs of holomorphic sections of  $L_\lambda$ .

*Remark 1.* (1) For  $k \in \mathbb{R}_{>0}$  and  $\lambda \in \mathfrak{t}_+^*$ ,  $\mathcal{O}_{k\lambda} = \mathcal{O}_\lambda$  as complex manifolds. If we compare the symplectic forms and the moment maps under this identification, we have  $\omega_{k\lambda} = k\omega_\lambda$  and  $\iota_{k\lambda} = k\iota_\lambda$ . In the case  $k \in \mathbb{Z}_{>0}$  and  $\lambda \in P_+$ , we have  $L_{k\lambda} \cong L_\lambda^{\otimes k}$ .

(2) For  $\lambda \in P_+$ , the action on  $\mathcal{O}_\lambda$  of the center  $Z(G) \cong \mathbb{Z}/3\mathbb{Z}$  of  $G = SU(3)$  is trivial, while that on  $L_\lambda$  is not trivial in general. However, by the construction of  $L_\lambda$  (see, e.g., [12]), if we suppose  $\lambda \in P_+ \cap Q$ , this action becomes trivial, too.

### § 2.3. Double weight varieties

For general properties of symplectic and Kähler quotients, see, e.g., [10], [14] and [18]. The following still holds for a general compact Lie group  $G$ .

Let  $\lambda_1, \lambda_2 \in \mathfrak{t}_+^*$ . The diagonal action of the maximal torus  $T$  of  $G$  on the direct product  $\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}$  is also Hamiltonian and the moment map  $\Phi : \mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2} \rightarrow \mathfrak{t}^*$  is

given by the composition of the map

$$\begin{aligned}\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2} &\rightarrow \mathfrak{g}^* \\ (x_1, x_2) &\rightarrow x_1 + x_2\end{aligned}$$

and the projection  $\mathfrak{g}^* \rightarrow \mathfrak{t}^*$ . For  $\mu \in \mathfrak{t}^*$ , we define the symplectic (or Kähler) quotient at  $\mu$  by

$$(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T := \Phi^{-1}(\mu) / T.$$

Here we assume that for  $\mu \in \mathfrak{t}^*$ ,

$$(a0) \quad \Phi^{-1}(\mu) \neq \emptyset,$$

$$(a1) \quad \mu \text{ is a regular value of the moment map } \Phi, \text{ and } (\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T \text{ is a smooth manifold.}$$

Then there exist a natural symplectic structure  $\omega = \omega(\lambda_1, \lambda_2, \mu)$  and a compatible complex structure  $J$  on  $(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T$ , induced from those on  $\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}$ , which make  $(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T$  a Kähler manifold. The complex dimension  $d$  of  $(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T$  is

$$d = \dim_{\mathbb{R}} G - \frac{1}{2}(\dim_{\mathbb{R}} G_{\lambda_1} + \dim_{\mathbb{R}} G_{\lambda_2}) - \dim_{\mathbb{R}} T.$$

We call  $(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T$  a *double weight variety*.

Now, suppose  $\lambda_1, \lambda_2 \in P_+$ . Let  $L_{\lambda_i}$  be the  $T$ -equivariant holomorphic line bundle over  $\mathcal{O}_{\lambda_i}$  as in Section 2.2, and let us set

$$\begin{aligned}\mathcal{L} &= (L_{\lambda_1} \boxtimes L_{\lambda_2}) //_{\mu} T := ((\mathrm{pr}_1^* L_{\lambda_1} \otimes \mathrm{pr}_2^* L_{\lambda_2})|_{\Phi^{-1}(\mu)}) / T, \\ \mathcal{L}_i &= L_{\lambda_i} //_{\mu} T := (\mathrm{pr}_i^* L_{\lambda_i}|_{\Phi^{-1}(\mu)}) / T \quad (i = 1, 2),\end{aligned}$$

where

$$\mathrm{pr}_i : \mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2} \rightarrow \mathcal{O}_{\lambda_i} \quad (i = 1, 2)$$

is the  $i$ -th projection.

By the assumption (a1), the isotropy subgroup at each point in  $\Phi^{-1}(\mu)$  is a finite group. Since its action on  $\mathrm{pr}_i^* L_{\lambda_i}$  ( $i = 1, 2$ ) and  $L_{\lambda_1} \boxtimes L_{\lambda_2} = \mathrm{pr}_1^* L_{\lambda_1} \otimes \mathrm{pr}_2^* L_{\lambda_2}$  may not be trivial,  $\mathcal{L}$  and  $\mathcal{L}_i$  ( $i = 1, 2$ ) are, in general, orbifold holomorphic line bundles over  $(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T$ . We assume that

$$(a2) \quad \mathcal{L}_i \text{ is a genuine holomorphic line bundle over } (\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T \text{ for } i = 1, 2.$$

Then we have  $\mathcal{L} = \mathcal{L}_1 \otimes \mathcal{L}_2$  and  $c_1(\mathcal{L}) = c_1(\mathcal{L}_1) + c_1(\mathcal{L}_2) = [\omega(\lambda_1, \lambda_2, \mu)]$ .

In general, let  $(M, \omega)$  be a compact Kähler manifold and suppose that there exists a holomorphic line bundle  $L$  over  $M$  such that  $c_1(L) = [\omega]$ . In this case, we define

$$\chi(M, L) := \sum_i (-1)^i H^i(M, \mathcal{O}(L))$$

as a virtual vector space. The dimension  $\dim_{\mathbb{C}} \chi(M, L)$  is called the Riemann-Roch number of  $(M, L)$ . We define the symplectic volume  $\text{vol}(M)$  of  $(M, \omega)$  by

$$\text{vol}(M) := \int_M \frac{\omega^d}{d!},$$

where  $d$  is the complex dimension of  $M$ .

**Lemma 2.1.** *Suppose that  $(M, \omega)$  and  $L$  satisfy the assumptions above. Then we have*

$$\text{vol}(M) = \lim_{k \rightarrow \infty} \frac{1}{k^d} \cdot \dim_{\mathbb{C}} \chi(M, L^{\otimes k}).$$

*Proof.* By the Hirzebruch-Riemann-Roch theorem, we have

$$\dim_{\mathbb{C}} \chi(M, L) = \int_M \text{Ch}(L) \text{Td}(M) = \int_M e^{c_1(L)} \text{Td}(M),$$

where  $\text{Ch}(L)$  is the Chern character of  $L$  and  $\text{Td}(M)$  is the Todd class of  $M$ . Hence we have

$$\lim_{k \rightarrow \infty} \frac{1}{k^d} \cdot \dim_{\mathbb{C}} \chi(M, L^{\otimes k}) = \lim_{k \rightarrow \infty} \int_M \frac{e^{kc_1(L)}}{k^d} \text{Td}(M) = \int_M \frac{c_1(L)^d}{d!} = \text{vol}(M).$$

□

In addition, we assume that a torus  $T$  acts holomorphically on  $M$ , and this action lifts to  $L \rightarrow M$ . Then the action of  $T$  on  $M$  is Hamiltonian, and the moment map  $\Phi : M \rightarrow \mathfrak{t}^*$  is determined. For suitably chosen  $\mu \in P$ , we obtain the Kähler manifold  $M//_{\mu}T = \Phi^{-1}(\mu)/T$ , the line bundle  $L//_{\mu}T = (L|_{\Phi^{-1}(\mu)})/T$  over  $M//_{\mu}T$  and  $\chi(M//_{\mu}T, L//_{\mu}T)$  (see, e.g., [8] and [17]).

*Remark 2.* (1) It follows from the multiplicative property of  $\chi$  (see the appendix in [11]) that  $\chi(\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}, L_{\lambda_1} \boxtimes L_{\lambda_2}) = \chi(\mathcal{O}_{\lambda_1}, L_{\lambda_1}) \otimes \chi(\mathcal{O}_{\lambda_2}, L_{\lambda_2})$ .

(2) As in Remark 1 (1),  $(\mathcal{O}_{k\lambda_1} \times \mathcal{O}_{k\lambda_2})//_{k\mu}T = (\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T$  as complex manifolds, and  $\omega(k\lambda_1, k\lambda_2, k\mu) = k\omega(\lambda_1, \lambda_2, \mu)$  for  $k \in \mathbb{R}_{>0}$ ,  $\lambda_1, \lambda_2 \in \mathfrak{t}_+^*$  and  $\mu \in \mathfrak{t}^*$ . In particular, it follows that  $\text{vol}((\mathcal{O}_{k\lambda_1} \times \mathcal{O}_{k\lambda_2})//_{k\mu}T) = k^d \cdot \text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T)$ . In the case  $k \in \mathbb{Z}_{>0}$ ,  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$ , we have  $(L_{k\lambda_1} \boxtimes L_{k\lambda_2})//_{k\mu}T = ((L_{\lambda_1} \boxtimes L_{\lambda_2})//_{\mu}T)^{\otimes k}$ .

(3) Even if  $\lambda_1, \lambda_2 \in P_+$  does not satisfy (a2),  $n\lambda_1$  and  $n\lambda_2$  does satisfy (a2) for some positive integer  $n$ . Hence, as far as the symplectic volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T)$  is concerned, we can assume (a2) without loss of generality.

On the other hand, we can regard  $\chi(M, L)$  as a representation of  $T$ . Then we obtain the weight decomposition

$$\chi(M, L) = \sum_{\mu \in P} m_{\mu} W_{\mu},$$



where  $W_\mu$  denotes the weight space associated with  $\mu \in P$  and

$$m_\mu = [\chi(M, L) : W_\mu] := \dim_{\mathbb{C}}(\chi(M, L) \otimes W_\mu^*)^T$$

is the weight multiplicity of  $W_\mu$ . Besides, for a representation  $V$  of  $T$ ,

$$V^T := \{v \in V \mid t \cdot v = v \ (\forall t \in T)\}$$

is the set of invariants in  $V$ .

The Guillemin-Sternberg theorem and its generalization (see, e.g., [8] and [17]) tell us the following.

**Theorem 2.2.** *Assume that the action of  $T$  on  $\Phi^{-1}(\mu)$  is free. Then we have*

$$\dim_{\mathbb{C}} \chi(M//_{\mu}T, L//_{\mu}T) = \dim_{\mathbb{C}}(\chi(M, L) \otimes W_\mu^*)^T = [\chi(M, L) : W_\mu].$$

In our situation, we obtain the following.

**Proposition 2.3.** *Suppose that  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$  satisfy the assumptions (a0), (a1) and (a2). Then we have*

- (1)  $\dim_{\mathbb{C}} \chi((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T, (L_{\lambda_1} \boxtimes L_{\lambda_2})//_{\mu}T) = [V_{\lambda_1} \otimes V_{\lambda_2} : W_\mu],$
- (2)  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T) = \lim_{k \rightarrow \infty} \frac{1}{k^d} \cdot [V_{k\lambda_1} \otimes V_{k\lambda_2} : W_{k\mu}].$

*Proof.* (1) It follows from Theorem 2.2, Remark 2 (1), and the Borel-Weil theorem.

(2) By Lemma 2.1, Remark 2 (2), and the assertion (1), we obtain

$$\begin{aligned} & \text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T) \\ &= \lim_{k \rightarrow \infty} \frac{1}{k^d} \cdot \dim_{\mathbb{C}} \chi((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T, ((L_{\lambda_1} \boxtimes L_{\lambda_2})//_{\mu}T)^{\otimes k}) \\ &= \lim_{k \rightarrow \infty} \frac{1}{k^d} \cdot \dim_{\mathbb{C}} \chi((\mathcal{O}_{k\lambda_1} \times \mathcal{O}_{k\lambda_2})//_{k\mu}T, (L_{k\lambda_1} \boxtimes L_{k\lambda_2})//_{k\mu}T) \\ &= \lim_{k \rightarrow \infty} \frac{1}{k^d} \cdot [V_{k\lambda_1} \otimes V_{k\lambda_2} : W_{k\mu}]. \end{aligned}$$

□

In the next section, we will compute the Rimann-Roch number  $\dim_{\mathbb{C}} \chi((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T, (L_{\lambda_1} \boxtimes L_{\lambda_2})//_{\mu}T)$  and the volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu}T)$  for  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$ .

### § 3. Main results

#### § 3.1. Combinatorial expression

For  $\lambda \in P_+$ , let us set

$$\chi_\lambda := \frac{\sum_{w \in W} \varepsilon(w) e^{w(\lambda + \rho)}}{e^\rho \prod_{\alpha \in \Delta_+} (1 - e^{-\alpha})} \in \mathbb{C}[e^{\Lambda_1}, e^{\Lambda_2}][[e^{-\Lambda_1}, e^{-\Lambda_2}]]$$

and define

$$F_{\lambda_1, \lambda_2}^\mu := \chi_{\lambda_1} \cdot \chi_{\lambda_2} \cdot e^{-\mu}$$

for  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$ . As in the Weyl character formula, we also regard them as functions on  $T$ .

**Lemma 3.1.** *For  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$ ,  $[V_{\lambda_1} \otimes V_{\lambda_2} : W_\mu]$  is equal to the coefficient of  $e^0$  (i.e., the constant term) in  $F_{\lambda_1, \lambda_2}^\mu$ .*

*Proof.* Let  $d\mu_T$  be the normalized invariant measure on  $T$ . Then we have

$$(3.1) \quad [V_{\lambda_1} \otimes V_{\lambda_2} : W_\mu] = \int_T \chi_{\lambda_1}(t) \chi_{\lambda_2}(t) e^{-\mu}(t) d\mu_T.$$

Let  $H_1, H_2$  be the basis of the integral lattice in  $\mathfrak{t}$  as (2.1). we write an element  $t \in T$  as  $t = \exp(x_1 H_1 + x_2 H_2)$  with  $x_1, x_2 \in [0, 1]$ . Let us set  $u_i = e^{2\pi\sqrt{-1}x_i}$  ( $i = 1, 2$ ) and define an isomorphism  $T \cong U(1)^2$  by  $t \mapsto (u_1, u_2)$ . Then we have

$$d\mu_T = dx_1 dx_2 = \frac{du_1}{2\pi\sqrt{-1}u_1} \frac{du_2}{2\pi\sqrt{-1}u_2}.$$

Hence (3.1) is equal to the coefficient of  $u_1^0 u_2^0$  in  $F_{\lambda_1, \lambda_2}^\mu(t)$ , which is regarded as a Laurent series of  $(u_1, u_2)$ . If we write  $F_{\lambda_1, \lambda_2}^\mu$  as

$$F_{\lambda_1, \lambda_2}^\mu = \sum C_{m_1, m_2} e^{m_1 \Lambda_1 + m_2 \Lambda_2},$$

then we have

$$F_{\lambda_1, \lambda_2}^\mu(u_1, u_2) = \sum C_{m_1, m_2} u_1^{m_1} u_2^{m_2}.$$

Therefore, (3.1) is equal to the coefficient of  $e^0$  in  $F_{\lambda_1, \lambda_2}^\mu$ .  $\square$

**Proposition 3.2.** *Let  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$ . Then we have*

$$[V_{\lambda_1} \otimes V_{\lambda_2} : W_\mu] = \sum_{w_1, w_2 \in W} \varepsilon(w_1) \varepsilon(w_2) E(\lambda_1, \lambda_2, \mu; w_1, w_2),$$

where for  $w_1, w_2 \in W$ , we define

$$E(\lambda_1, \lambda_2, \mu; w_1, w_2) = \sum_{(j_1, j_2, j_3)} (j_1 + 1)(j_2 + 1)(j_3 + 1)$$

and the sum is taken over all  $(j_1, j_2, j_3) \in (\mathbb{Z}_{\geq 0})^3$  which satisfy the condition

$$(3.2) \quad w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - 2\rho - j_1\alpha_1 - j_2\alpha_2 - j_3(\alpha_1 + \alpha_2) = 0.$$

*Proof.* Applying the generalized binomial theorem to  $(e^\rho \prod_{\alpha \in \Delta_+} (1 - e^{-\alpha}))^{-2}$  in  $F_{\lambda_1, \lambda_2}^\mu$ , we have a power series expansion

$$\begin{aligned}
F_{\lambda_1, \lambda_2}^\mu &= \sum_{w_1, w_2 \in W} \sum_{(j_1, j_2, j_3)} \varepsilon(w_1) \varepsilon(w_2) (-1)^{j_1 + j_2 + j_3} \binom{-2}{j_1} \binom{-2}{j_2} \binom{-2}{j_3} \\
&\quad \cdot e^{w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - j_1 \alpha_1 - j_2 \alpha_2 - j_3(\alpha_1 + \alpha_2) - 2\rho} \\
&= \sum_{w_1, w_2 \in W} \sum_{(j_1, j_2, j_3)} \varepsilon(w_1) \varepsilon(w_2) \binom{j_1 + 1}{1} \binom{j_2 + 1}{1} \binom{j_3 + 1}{1} \\
&\quad \cdot e^{w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - j_1 \alpha_1 - j_2 \alpha_2 - j_3(\alpha_1 + \alpha_2) - 2\rho} \\
&= \sum_{w_1, w_2 \in W} \sum_{(j_1, j_2, j_3)} \varepsilon(w_1) \varepsilon(w_2) (j_1 + 1)(j_2 + 1)(j_3 + 1) \\
&\quad \cdot e^{w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - j_1 \alpha_1 - j_2 \alpha_2 - j_3(\alpha_1 + \alpha_2) - 2\rho}.
\end{aligned}$$

Hence our claim follows from Proposition 3.1.  $\square$

### § 3.2. Formulas for the Rimann-Roch number and the volume

In the following, we assume that  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$  satisfy the following three conditions.

(A0)  $\mu \in \text{conv.}\{w_1 \lambda_1 + w_2 \lambda_2 | w_1, w_2 \in W\}$ ,

(A1) For any  $w_1, w_2 \in W$ ,

$$\begin{aligned}
&\langle w_1 \lambda_1 + w_2 \lambda_2 - \mu, \Lambda_1 \rangle \neq 0, \\
&\langle w_1 \lambda_1 + w_2 \lambda_2 - \mu, \Lambda_2 \rangle \neq 0, \\
&\langle w_1 \lambda_1 + w_2 \lambda_2 - \mu, \Lambda_1 \rangle \neq \langle w_1 \lambda_1 + w_2 \lambda_2 - \mu, \Lambda_2 \rangle,
\end{aligned}$$

(A2)  $\lambda_1, \lambda_2, \mu \in Q$ ,

where  $\text{conv.}\{w_1 \lambda_1 + w_2 \lambda_2 | w_1, w_2 \in W\}$  denotes the convex full of  $\{w_1 \lambda_1 + w_2 \lambda_2 | w_1, w_2 \in W\}$ .

*Remark 3.* We mention the relation between the conditions (a0)–(a2) and the conditions (A0)–(A2). By the convexity theorem of Hamiltonian torus actions on symplectic manifolds (see, e.g., [1], [6] and [7]), (A0) implies (a0), and (A1) implies the former part of (a1). Moreover, for  $G = SU(3)$ , the former part of (a1) implies the latter part of (a1). Besides, by Remark 1 (2), (A2) implies (a2). As we noted in Remark 2 (3), as far as the symplectic volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T)$  is concerned, we may assume (A2) without loss of generality.

Under the conditions above, let us concretely compute  $E(\lambda_1, \lambda_2, \mu; w_1, w_2)$  in Proposition 3.2. The condition (3.2) in Proposition 3.2 means that  $(j_1, j_2, j_3) \in \mathbb{Z}^3$  satisfies

$$\begin{cases} j_1 = \langle w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - 2\rho, \Lambda_1 \rangle - j_3 \geq 0, \\ j_2 = \langle w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - 2\rho, \Lambda_2 \rangle - j_3 \geq 0, \\ j_3 \geq 0. \end{cases}$$

We write

$$\begin{aligned} A &= \langle w_1\lambda_1 + w_2\lambda_2 - \mu, \Lambda_1 \rangle, \quad B = \langle w_1\lambda_1 + w_2\lambda_2 - \mu, \Lambda_2 \rangle, \\ C &= \langle w_1\rho + w_2\rho, \Lambda_1 \rangle, \quad D = \langle w_1\rho + w_2\rho, \Lambda_2 \rangle \end{aligned}$$

for brevity. Note that  $\langle \rho, \Lambda_1 \rangle = \langle \rho, \Lambda_2 \rangle = 1$ . Then the condition (3.2) means that  $(j_1, j_2, j_3) \in \mathbb{Z}^3$  satisfies

$$\begin{cases} j_1 = A + C - 2 - j_3 \geq 0, \\ j_2 = B + D - 2 - j_3 \geq 0, \\ j_3 \geq 0. \end{cases}$$

Therefore we consider the following cases. Recall that

$$\bar{\gamma}_1 = \mathbb{R}_{\geq 0}\alpha_1 + \mathbb{R}_{\geq 0}(\alpha_1 + \alpha_2), \quad \bar{\gamma}_2 = \mathbb{R}_{\geq 0}(\alpha_1 + \alpha_2) + \mathbb{R}_{\geq 0}\alpha_2.$$

Case 1. Suppose that  $A + C - 2 \geq 0$  and  $B + D - 2 \geq 0$ .

(1-1) If  $A + C \geq B + D$ , that is,  $w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - 2\rho \in \bar{\gamma}_1$ , then we have

$$\begin{aligned} E(\lambda_1, \lambda_2, \mu; w_1, w_2) &= \sum_{j_3=0}^{B+D-2} (A + C - 1 - j_3)(B + D - 1 - j_3)(j_3 + 1) \\ &= \frac{1}{12}(B + D - 1)(B + D)(B + D + 1)(2(A + C) - (B + D)). \end{aligned}$$

(1-2) If  $A + C \leq B + D$ , that is,  $w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - 2\rho \in \bar{\gamma}_2$ , then we have

$$\begin{aligned} E(\lambda_1, \lambda_2, \mu; w_1, w_2) &= \sum_{j_3=0}^{A+C-2} (A + C - 1 - j_3)(B + D - 1 - j_3)(j_3 + 1) \\ &= \frac{1}{12}(A + C - 1)(A + C)(A + C + 1)(-(A + C) + 2(B + D)). \end{aligned}$$

Case 2. Suppose that  $A + C - 2 < 0$  or  $B + D - 2 < 0$ .

We have  $E(\lambda_1, \lambda_2, \mu; w_1, w_2) = 0$ .

Combining all the results above, we obtain the following explicit formula for the Rimann-Roch number  $\dim_{\mathbb{C}} \chi((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu} T, (L_{\lambda_1} \boxtimes L_{\lambda_2})//_{\mu} T)$ .

**Proposition 3.3.** *Let  $\lambda_1, \lambda_2 \in P_+$  and  $\mu \in P$  satisfy the assumptions (A0), (A1) and (A2). Then we have*

$$\dim_{\mathbb{C}} \chi((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu} T, (L_{\lambda_1} \boxtimes L_{\lambda_2})//_{\mu} T) = \sum_{w_1, w_2 \in W} \varepsilon(w_1) \varepsilon(w_2) E(\lambda_1, \lambda_2, \mu; w_1, w_2),$$

where  $E(\lambda_1, \lambda_2, \mu; w_1, w_2)$  is given as follows. We write

$$\begin{aligned} A &= \langle w_1 \lambda_1 + w_2 \lambda_2 - \mu, \Lambda_1 \rangle, \quad B = \langle w_1 \lambda_1 + w_2 \lambda_2 - \mu, \Lambda_2 \rangle, \\ C &= \langle w_1 \rho + w_2 \rho, \Lambda_1 \rangle, \quad D = \langle w_1 \rho + w_2 \rho, \Lambda_2 \rangle \end{aligned}$$

for brevity.

(1) If  $w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - 2\rho \in \bar{\gamma}_1$ ,

$$\begin{aligned} E(\lambda_1, \lambda_2, \mu; w_1, w_2) \\ = \frac{1}{12} (B + D - 1)(B + D)(B + D + 1)(2(A + C) - (B + D)). \end{aligned}$$

(2) If  $w_1(\lambda_1 + \rho) + w_2(\lambda_2 + \rho) - \mu - 2\rho \in \bar{\gamma}_2$ ,

$$\begin{aligned} E(\lambda_1, \lambda_2, \mu; w_1, w_2) \\ = \frac{1}{12} (A + C - 1)(A + C)(A + C + 1)(-(A + C) + 2(B + D)). \end{aligned}$$

(3) Otherwise,

$$E(\lambda_1, \lambda_2, \mu; w_1, w_2) = 0.$$

In the following, we assume that both of two coadjoint orbits  $\mathcal{O}_{\lambda_1}$  and  $\mathcal{O}_{\lambda_2}$  are flag manifolds of  $SU(3)$ . Namely, we assume  $\lambda_1, \lambda_2 \in P_{++}$ . Combining Propositions 2.3 and 3.3, we obtain the following formula for the volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu} T)$ .

**Theorem 3.4.** *Let  $\lambda_1, \lambda_2 \in P_{++}$  and  $\mu \in P$  satisfy the assumptions (A0), (A1) and (A2). Then we have*

$$(3.3) \quad \text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu} T) = \sum_{w_1, w_2 \in W} \varepsilon(w_1) \varepsilon(w_2) F(\lambda_1, \lambda_2, \mu; w_1, w_2),$$

where  $F(\lambda_1, \lambda_2, \mu; w_1, w_2)$  is given as follows.

(1) If  $w_1\lambda_1 + w_2\lambda_2 - \mu \in \gamma_1$ ,

$$F(\lambda_1, \lambda_2, \mu; w_1, w_2) = \frac{1}{12} \langle w_1\lambda_1 + w_2\lambda_2 - \mu, \Lambda_2 \rangle^3 \langle w_1\lambda_1 + w_2\lambda_2 - \mu, 2\Lambda_1 - \Lambda_2 \rangle.$$

(2) If  $w_1\lambda_1 + w_2\lambda_2 - \mu \in \gamma_2$ ,

$$F(\lambda_1, \lambda_2, \mu; w_1, w_2) = \frac{1}{12} \langle w_1\lambda_1 + w_2\lambda_2 - \mu, \Lambda_1 \rangle^3 \langle w_1\lambda_1 + w_2\lambda_2 - \mu, -\Lambda_1 + 2\Lambda_2 \rangle.$$

(3) Otherwise,

$$F(\lambda_1, \lambda_2, \mu; w_1, w_2) = 0.$$

*Proof.* We first note that for  $\lambda_1, \lambda_2 \in P_{++}$  and  $\mu \in P$  satisfying the assumptions (A0), (A1) and (A2), we have

$$d = \dim_{\mathbb{C}}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2}) //_{\mu} T) = 4.$$

According to Propositions 2.3 and 3.3, we must compute  $E(k\lambda_1, k\lambda_2, k\mu; w_1, w_2)$ . In this case, the condition in (1) of Proposition 3.3 means that  $kB + D - 2 \geq 0$  and

$$\begin{aligned} 0 &\leq \langle w_1(k\lambda_1 + \rho) + w_2(k\lambda_2 + \rho) - 2\rho - k\mu, \Lambda_1 - \Lambda_2 \rangle \\ &= k(A - B) + (C - D). \end{aligned}$$

Let us take  $k$  large enough. By the assumption (A1), these inequalities above means that  $A > B > 0$ , that is,  $w_1\lambda_1 + w_2\lambda_2 - \mu \in \gamma_1$ . Then we have

$$\begin{aligned} E(k\lambda_1, k\lambda_2, k\mu; w_1, w_2) &= \frac{1}{12} (kB + D - 1)(kB + D)(kB + D + 1)(2(kA + C) - (kB + D)) \\ &= \frac{k^4}{12} B^3(2A - B) + (\text{lower terms of } k). \end{aligned}$$

Similarly, when  $k \gg 0$ , the condition in (2) of Proposition 3.3 means that  $B > A > 0$ , that is,  $w_1\lambda_1 + w_2\lambda_2 - \mu \in \gamma_2$ . Then we have

$$\begin{aligned} E(k\lambda_1, k\lambda_2, k\mu; w_1, w_2) &= \frac{1}{12} (kA + C - 1)(kA + C)(kA + C + 1)(-(kA + C) + 2(kB + D)) \\ &= \frac{k^4}{12} A^3(-A + 2B) + (\text{lower terms of } k). \end{aligned}$$

Finally, when  $k \gg 0$ , the condition in (3) of Proposition 3.3 means that  $A < 0$  or  $B < 0$ . Then we have

$$E(k\lambda_1, k\lambda_2, k\mu; w_1, w_2) = 0.$$

Combining all the results above and Proposition 2.3, we obtain our claim.  $\square$

*Remark 4.* (1) In the case where  $\lambda_1$  or  $\lambda_2$  is in  $P_+ - P_{++}$ , although  $F(\lambda_1, \lambda_2, \mu; w_1, w_2) \neq 0$  in general, the right hand side of the formula (3.3) is always equal to 0. We need another consideration for this. We will discuss it in the forthcoming paper [19].

(2) In general, for an element  $\mu$  in the weight lattice  $P$ , there exist  $w \in W$  and  $\mu' \in P_+$  such that  $\mu = w\mu'$ . Since  $[V_{\lambda_1} \otimes V_{\lambda_2} : W_\mu] = [V_{\lambda_1} \otimes V_{\lambda_2} : W_{\mu'}]$ , we have  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_\mu T) = \text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_{\mu'} T)$ . But the computation of the right hand side of the formula (3.3) for  $\mu' \in P_+$  becomes more simpler than that for  $\mu \in P$ .

Now, let  $\lambda_1, \lambda_2 \in (P \otimes \mathbb{Q}) \cap \mathfrak{t}_{++}^*$  and  $\mu \in P \otimes \mathbb{Q}$  satisfy the conditions (A0) and (A1). Then there exists  $n \in \mathbb{Z}_{>0}$  such that  $n\lambda_1, n\lambda_2 \in P_{++}$  and  $n\mu \in P$ . Therefore it follows from Remark 2 (2) and Proposition 2.3 that the volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_\mu T)$  is given by the right hand side of the formula (3.3). Furthermore, by continuity of the symplectic volume of symplectic quotients (see, e.g., [3] and [9]), we obtain the following.

**Corollary 3.5.** *Suppose that  $\lambda_1, \lambda_2 \in \mathfrak{t}_{++}^*$  and  $\mu \in \mathfrak{t}^*$  satisfy the conditions (A0) and (A1) in Section 3.2. Then the volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_\mu T)$  is given by the right hand side of the formula (3.3).*

## § 4. An example

As a simple example of Theorem 3.4 or Corollary 3.5, let us consider the case where  $\mu$  is sufficiently close to  $\lambda_1 + \lambda_2$ . In this case, we have  $F(\lambda_1, \lambda_2, \mu; w_1, w_2) = 0$  unless the case  $w_1 = w_2 = e$ . Hence we have

$$\begin{aligned} \text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_\mu T) &= F(\lambda_1, \lambda_2, \mu; w_1 = e, w_2 = e) \\ &= \begin{cases} \frac{1}{12} \langle \lambda_1 + \lambda_2 - \mu, \Lambda_2 \rangle^3 \langle \lambda_1 + \lambda_2 - \mu, 2\Lambda_1 - \Lambda_2 \rangle & (\text{if } \lambda_1 + \lambda_2 - \mu \in \gamma_1), \\ \frac{1}{12} \langle \lambda_1 + \lambda_2 - \mu, \Lambda_1 \rangle^3 \langle \lambda_1 + \lambda_2 - \mu, -\Lambda_1 + 2\Lambda_2 \rangle & (\text{if } \lambda_1 + \lambda_2 - \mu \in \gamma_2). \end{cases} \end{aligned}$$

Now if we write  $\lambda_1 = p\alpha_1 + q\alpha_2, \lambda_2 = r\alpha_1 + s\alpha_2 \in P_{++}$  and  $\mu = x\alpha_1 + y\alpha_2 \in P$ , we can express the volume  $\text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_\mu T)$  as a polynomial of  $p, q, r, s, x, y$  as follows.

$$\begin{aligned} \text{vol}((\mathcal{O}_{\lambda_1} \times \mathcal{O}_{\lambda_2})//_\mu T) &= \begin{cases} \frac{1}{12} (q + s - y)^3 (2p - q + 2r - s - 2x + y) & (\text{if } \lambda_1 + \lambda_2 - \mu \in \gamma_1), \\ \frac{1}{12} (p + r - x)^3 (-p + 2q - r + 2s + x - 2y) & (\text{if } \lambda_1 + \lambda_2 - \mu \in \gamma_2). \end{cases} \end{aligned}$$

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